

A Mesoscopic Particle-based Method for Active-Nematics



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Physics and Biology of Active
Systems
23 June, 2015

Microscopic Models

- Active colloids
- Simulations of motile microbes

Motivation



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- Simulations of motile microbes

Continuum Models

- Active-nematic hydrodynamics



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Mescoscale Models

- Coarse-grain many microscopic entities without calculating molecular interactions

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Mescoscale Models

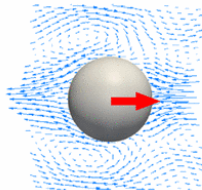
- Coarse-grain many microscopic entities without calculating molecular interactions

Continuum Models

- Active-nematic hydrodynamics

Conclusion:

- Coarse-graining solvents successful in passive soft matter
 - We would like a “toolbox” of similar simulation techniques for active fluids



Colloids in Active Fluids: Anomalous Microrheology and Negative Drag by Foffano, *et. al.*

Active-Nematic Hydrodynamic Equations to be Simulated

- Continuity equation $D_t \rho + \rho \partial_i u_i = 0$

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 - viscous stress Π_{ij}^{visc}
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 - force dipole density
 - ζ^{act} positive for extensile and negative for contractile

Active-Nematic

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- Review: recreate Navier-Stokes Equations by mesoscopic **MPCD**

Traditional MPCD

- Particle-based algorithm for moderate Péclet numbers.

Each point-like particle i has

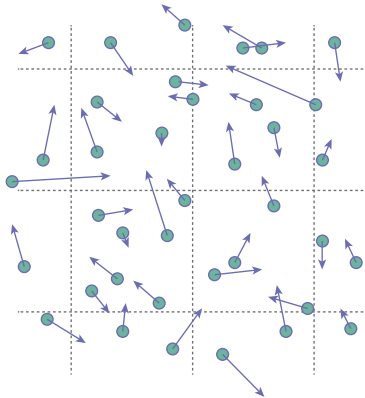
1. position $\underline{x}_i$
2. velocity $\underline{v}_i$
3. mass m_i

Each cell has

1. population N_c
2. centre of mass velocity $\underline{v}_{\text{cm}}$
3. (moment of inertia $\underline{I}_{\text{c}}$)

Advantages

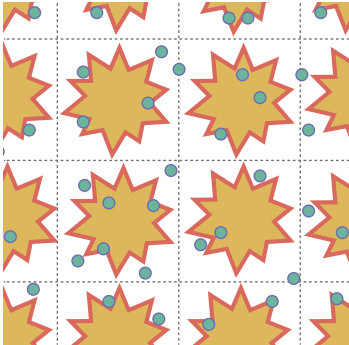
- Finite temperature
- No explicit pair-wise interactions
 - Point-like particles
- Highly parallelizable



Streaming

$$\underline{x}_i(t + \delta t) = \underline{x}_i(t) + \underline{v}_i(t) \delta t$$

† Modelling the separation of macromolecules: A review of current computer simulation methods, Slater *et. al.* , Electrophoresis, 30, 2009.



Streaming

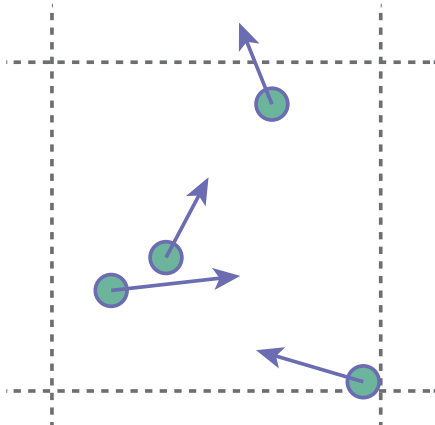
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Collision (Andersen MPCD)

$$\underline{v}_i(t + \delta t) = \underline{v}_{\text{CM}}(t) + \Xi$$

where Ξ is a collision operation

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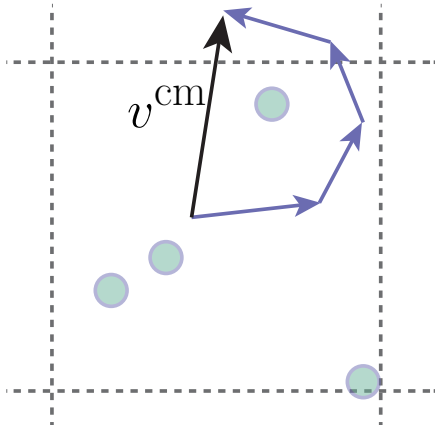
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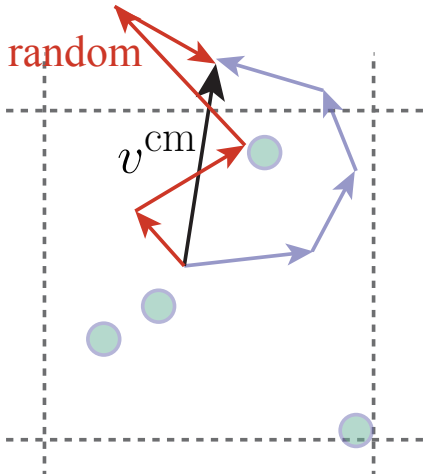
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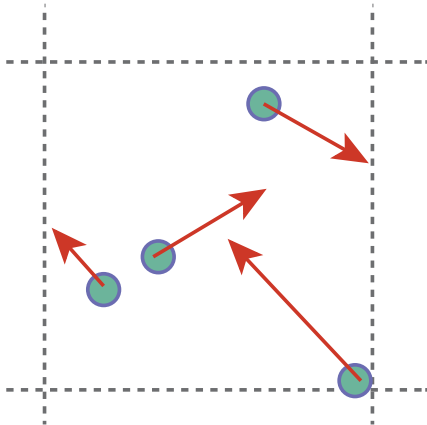
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So far...

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- Momentum equation $\rho D_t u_i = \partial_j [\Pi_{ij}^{\text{visc}} + \Pi_{ij}^{\text{nem}} + \Pi_{ij}^{\text{act}}]$
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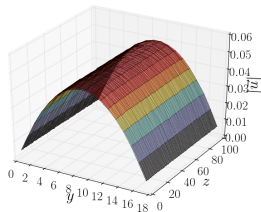
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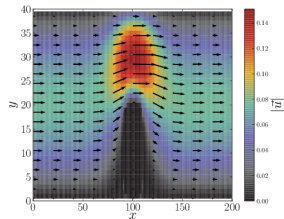
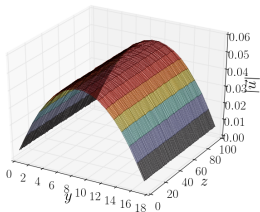
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- *i.e.* **Navier-Stokes Equations** on long time and length scales

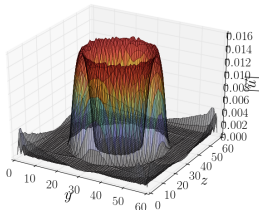
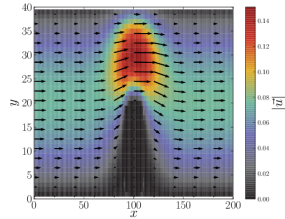
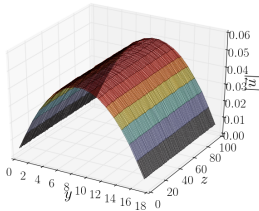
Example: MPCD Flow Profiles



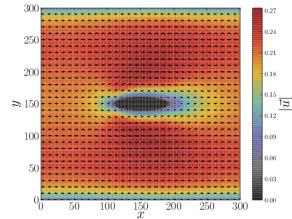
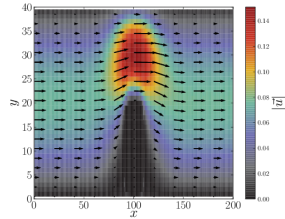
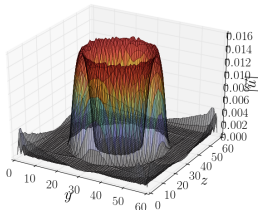
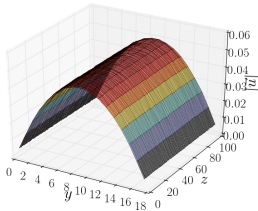
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Questions on the **operation** of the traditional MPCD algorithm?

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Nematic MPCD

Each must additionally have

1. orientation $\underline{u}_i$
2. molecular potential U
3. rotational coefficient γ_R
4. tumbling parameter λ

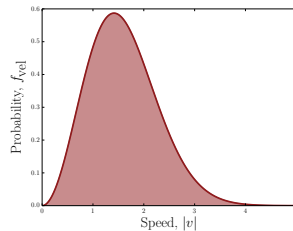
Each cell must have

1. order parameter $\underline{Q}_c$
 - 1.1 scalar order parameter S_c
 - 1.2 director $\underline{n}_c$

Momentum Collision Operation

$$\underline{v}_i(t + \delta t) = \underline{v}_{\text{CM}}(t) + \Xi$$

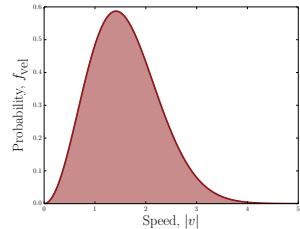
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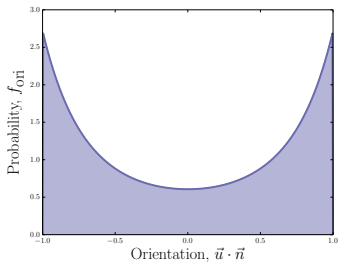
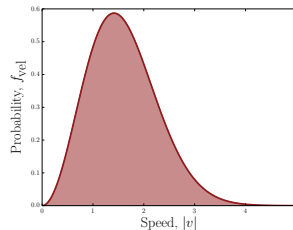
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Orientation Collision Operation

$$\underline{u}_i(t + \delta t) = \underline{\eta}_i(\beta U, S, \underline{n}),$$

where $\underline{\eta}_i$ is a random orientation drawn from the equilibrium **Maier-Saupe** probability distribution

$$f_{\text{ori}} = \exp\left(\frac{\beta U S}{d-1} \left[d(\underline{\eta} \cdot \underline{n})^2 - 1\right]\right) \propto \exp\left(\frac{d\beta U S}{d-1} \eta_n^2\right)$$

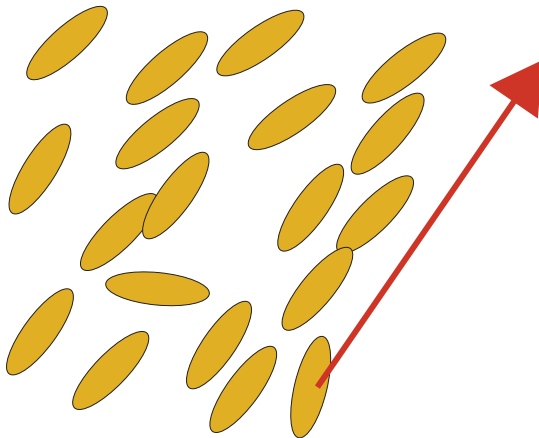


Figure: Nematic $S \approx 1$; $U \gg k_B T$

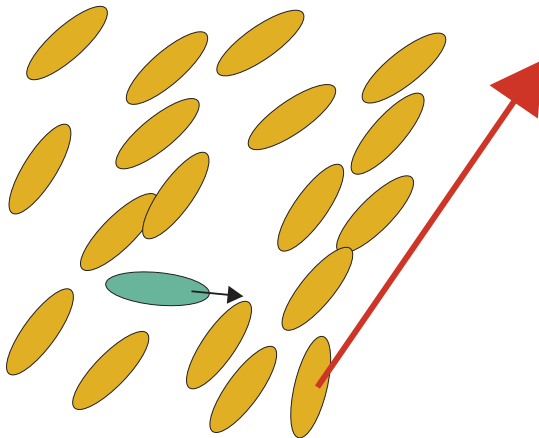


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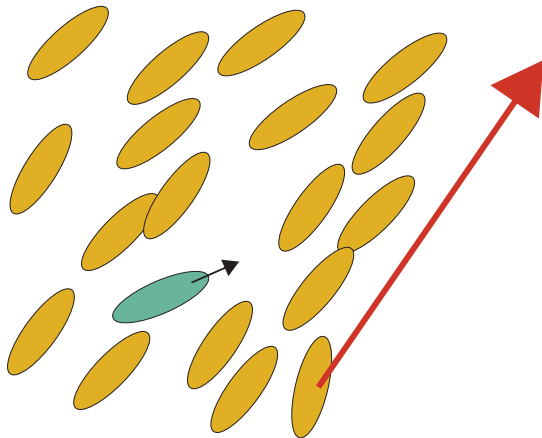
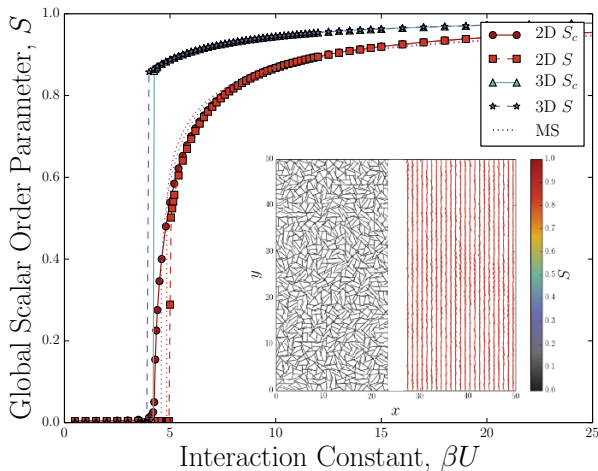
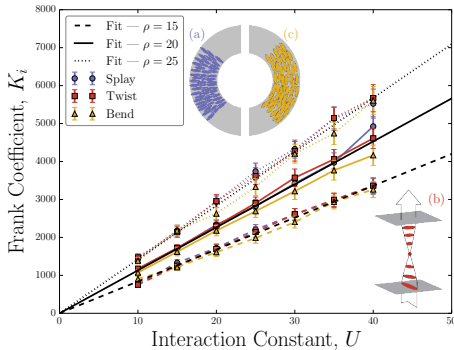


Figure: Nematic $S \approx 1$; $U \gg k_B T$

Figure: Isotropic phase ($U \ll k_B T$)

Figure: Nematic phase (note: topological defects)





Isotropic Elasticity

$$K_i = \ell^2 \rho U S^2 [1 + C_i] / 6 \sim U/a$$

- $i = \{\text{splay, twist, bend}\}$
- $C_i \approx 0$ (rod length small compared to interaction length)
- $\ell \approx a$
- isotropic elasticity and **one-constant approximation** applies.

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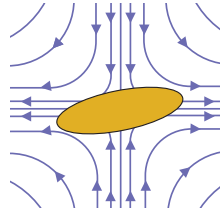
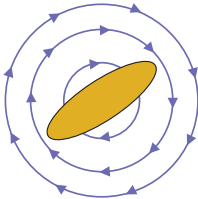
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- **Nematohydrodynamics** requires coupling momentum and orientation

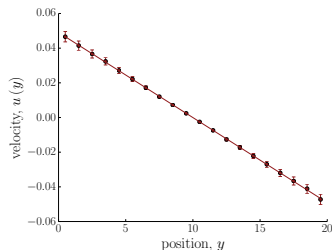
Shear Alignment: Velocity $\rightarrow$ Orientation Couplings



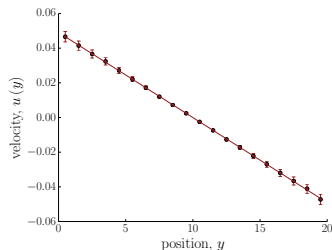
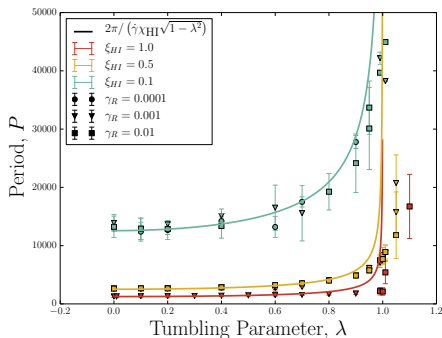
Flow-Induced Rotation

- Discretised Jeffery's equation for a slender rod

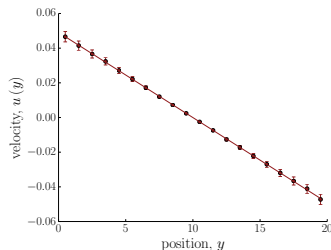
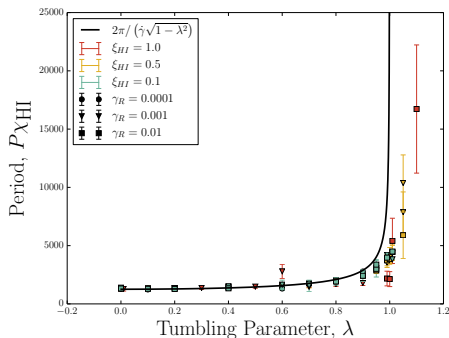
$$\frac{\delta \underline{u}_{HI,i}}{\delta t} = \chi_{HI} \left[\underline{u}_i \cdot \underline{\underline{\omega}} + \lambda \left(\underline{u}_i \cdot \underline{\underline{D}} - \underline{u}_i \underline{u}_i \underline{u}_i : \underline{\underline{D}} \right) \right]$$
- bare tumbling parameter λ
- heuristic shear coupling coefficient χ_{HI}



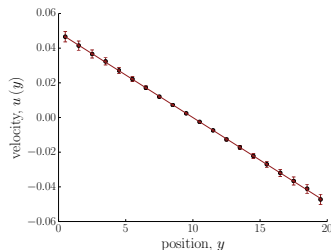
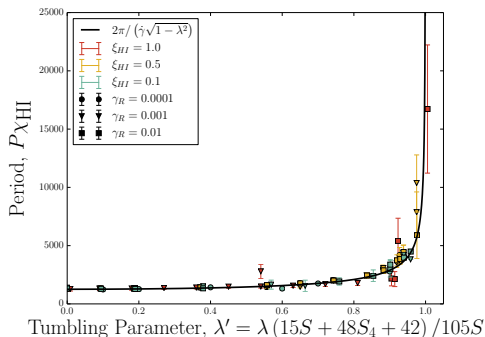
- Lees-Edwards BCs
- If $\lambda < 1$ then tumbling behaviour



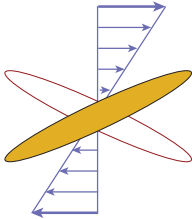
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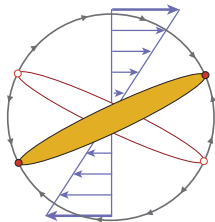


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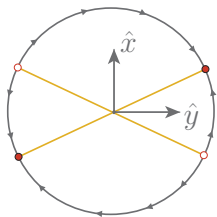
Expect:

- If $\lambda > 1$ then shear aligning



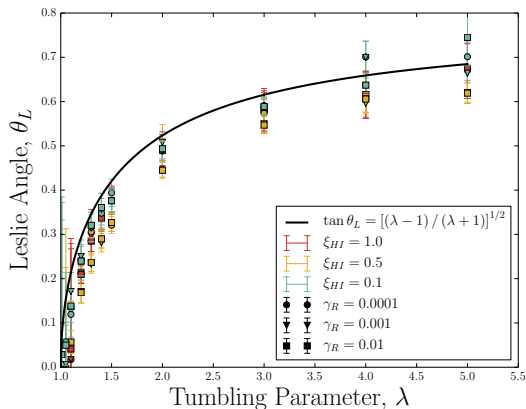
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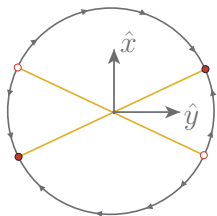
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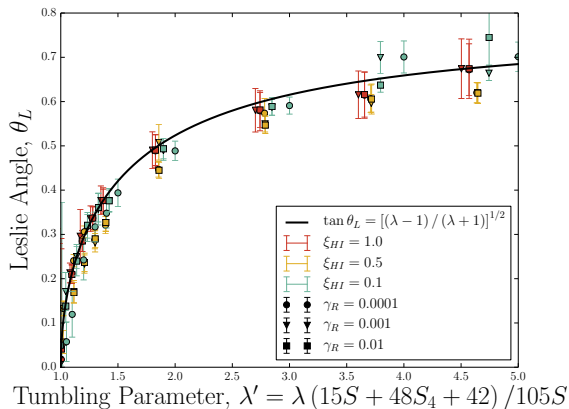
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 - Momentum equation $\rho D_t u_i = \partial_j [\Pi_{ij}^{\text{visc}} + \Pi_{ij}^{\text{nem}} + \Pi_{ij}^{\text{act}}]$
 - Orientation equation $D_t Q_{ij} - S_{ij} = \Gamma H_{ij}$
-
- Backflow next

Backflow: Orientation → Velocity Coupling

Torque Balance

$$\begin{aligned}\underline{\Gamma}_{\text{net},i} &= \underline{\Gamma}_{\text{HI},i} + \underline{\Gamma}_{\text{col},i} + \underline{\Gamma}_{\text{ext},i} = 0 \\ \underline{\Gamma}_{\text{HI},i} &= -\underline{\Gamma}_{\text{col},i} - \underline{\Gamma}_{\text{ext},i}\end{aligned}$$

Hydrodynamic Drag Passes Angular Momentum To Fluid

The MPCD collision operator modified to

$$\Xi_{i,c} = \underline{\xi}_i - \left\langle \underline{\xi}_j \right\rangle_{N_c} + \left(\underline{I}_c^{-1} \cdot [\delta \underline{L}_c + \delta \underline{\mathcal{L}}_c] \right) \times \underline{r}'_i$$

where $\delta \underline{\mathcal{L}}_c = \sum_i^{N_c} \delta \underline{\mathcal{L}}_i = - \sum_i^{N_c} \underline{\Gamma}_{\text{HI},i} \delta t$

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 - Remaining: **Activity**

Questions on the **operation** of the
nematic-MPCD algorithm?

Collision Operator

Passive Fluids

- MPCD collisions are **artificial operations** that **stochastically** exchange velocities while **conserving** energy and momentum

Active Fluids

- Active materials convert chemical potential energy to spontaneously drive dynamics
- MPCD collision operations can be conceived that lead to various classes of active matter

Dipole-Force Andersen Collision

- MPCD cells are subject to dipolar force

$$\Xi_{i,c} = \underbrace{\xi_i}_{\text{noise}} - \underbrace{\langle \xi_j \rangle}_{\text{av. noise}} + \underbrace{\left(\underline{I}_c^{-1} \cdot [\delta \underline{L}_c + \delta \underline{\mathcal{L}}_c] \right)}_{\text{ang. mom.}} \times \underline{r}'_i$$

$$+ \underbrace{\tau [\alpha_{\text{act}} - |v_i(t)|] (\kappa_i \underline{n}_c)}_{\text{activity}}$$

$$- \underbrace{\langle \tau [\alpha_{\text{act}} - |v_j(t)|] (\kappa_j \underline{n}_c) \rangle_{N_c}}_{\text{residuals}}$$

where

- τ relaxation coefficient
 - Extensile $0 < \tau \leq 1$
 - Contractile $-1 \leq \tau < 0$
- α_{act} out-of-equilibrium speed
- $\kappa_i = \pm 1$ based on **plane**:

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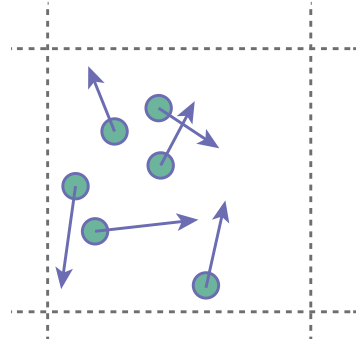
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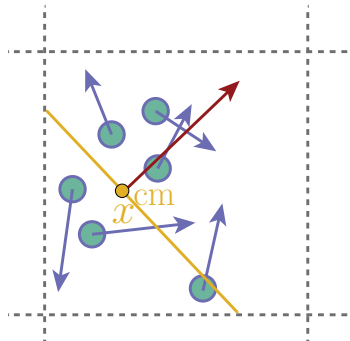
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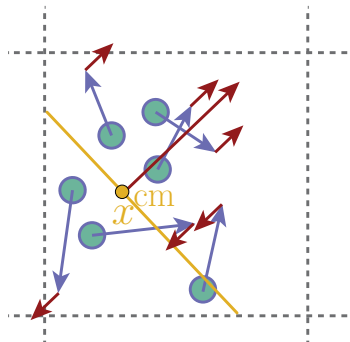
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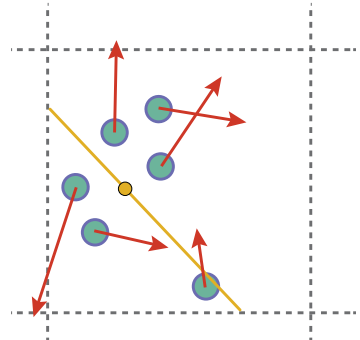
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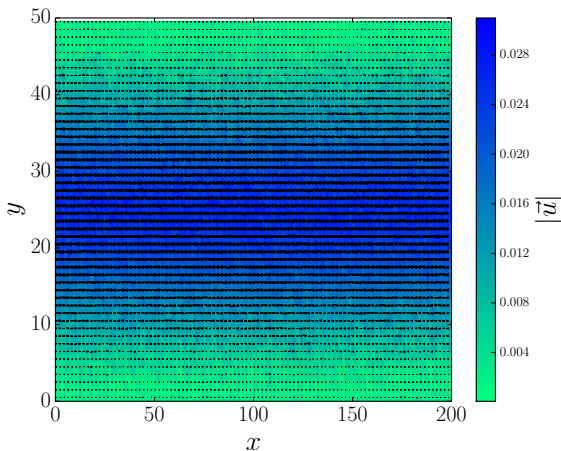
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Wall and Defect Pairs Formation



Wall and Defect Pairs Formation



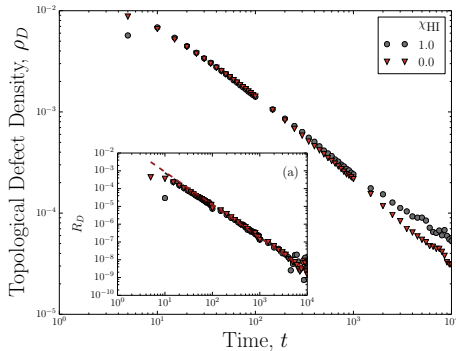


Active Colloidal Liquid Crystals



Recapitulation

- MPCD traditional **mesoscale** algorithm for isotropic solvent
- MPCD for nematic liquid crystals
 - Isotropic-nematic transition
 - Topological defect dynamics
 - Frank elastic coefficients
 - Nematodynamics with backflow
 - Tumbling and aligning
- MPCD for active nematic
 - Modified collision operator
 - Dipolar active fluid



Annihilation Rate

- $R_D = -\dot{\rho}_D \sim t^{-(\nu+1)}$
- Liu and Muthukumar: $\nu = 6/7$
- Varies in time: $\nu = 0.74 \pm 0.02$ for $t \in [10, 10^3]$ but $\nu = 0.83 \pm 0.04$ over $t \in [10^2, 10^4]$